

# **Magma ocean outgassing and hydrogen isotopic constraints on the Hadean Earth**

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5 Figures

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1 **Abstract**

2 The Moon-forming giant impact extensively melts and partially vaporizes the silicate  
3 Earth and delivers a substantial mass of metal to Earth's core. Subsequent evolution of  
4 the terrestrial magma ocean and overlying atmosphere has been described by theoretical  
5 models but observable constraints on this epoch have proved elusive. Here, we report  
6 calculations of primordial atmospheric evolution during the magma ocean and water  
7 ocean epochs and forge new links with observations to gain insight into the behavior of  
8 volatiles on the Hadean Earth. As accretion wanes, Earth's magma ocean crystallizes,  
9 outgassing the bulk of its volatiles into the primordial atmosphere. The redox state of the  
10 magma ocean controls both the chemical composition of the outgassed volatiles and the  
11 hydrogen isotopic composition of water oceans that remain after hydrogen escape from  
12 the primordial atmosphere. While water condenses and is retained, molecular hydrogen  
13 does not condense and can escape, allowing large quantities ( $\sim 10^2$  bars) of hydrogen to  
14 be, in principle, lost from Earth in this epoch. Because the escaping inventory of H can  
15 be comparable to the hydrogen inventory in primordial water oceans, the corresponding  
16 deuterium enrichment can be large with a magnitude that depends on the initial H<sub>2</sub>  
17 inventory of the atmosphere. Under equilibrium partitioning, the water molecule  
18 concentrates deuterium and, to the extent that hydrogen in other forms (e.g., H<sub>2</sub>) are  
19 significant species in the outgassed primordial atmosphere of Earth, pronounced D/H  
20 enrichments ( $\sim 2x$ ) in the oceans are expected to arise from equilibrium D/H partitioning  
21 alone. By contrast, the common view that terrestrial water has a primarily carbonaceous  
22 chondritic source requires seawater to preserve the isotopic composition of that source,  
23 undergoing minimal D-enrichment via equilibration with H<sub>2</sub> followed by hydrodynamic  
24 escape. Such minimal enrichment places upper limits on the amount of primordial  
25 atmospheric H<sub>2</sub> in contact with the early water oceans and implies relatively oxidizing

26 conditions ( $\log f_{\text{O}_2} > \text{IW} + 1$ ) for last equilibration with the magma ocean. Preservation of  
27 an approximate carbonaceous chondrite D/H signature in seawater thus provides  
28 evidence that the observed oxidation of silicate Earth occurred before or during the  
29 crystallization of the final magma ocean, yielding a new constraint on the timing of this  
30 critical event in Earth history. The seawater-carbonaceous chondrite “match” in D/H (to  
31  $\sim 10\text{-}20\%$ ) further constrains the prior existence of an atmospheric  $\text{H}_2$  inventory – of any  
32 origin – on the post-giant-impact Earth to  $< 20$  bars, and suggests that the terrestrial  
33 mantle supplied the oxidant for the chemical destruction of metals during terrestrial late  
34 accretion.

35 Keywords: magma ocean; outgassing; water; hydrogen; isotope; Hadean

36

## 37 **1. Introduction**

38 The composition and origin of Earth’s early atmosphere has been debated since at least  
39 the mid-twentieth century (Brown, 1949). Recent interest arises from a desire to  
40 understand climate on the early Earth (Wordsworth and Pierrehumbert, 2013a) as well  
41 as the environment that led to abiogenesis (Kasting, 2014), and because the volatile  
42 history of Earth gives insight into the origin and evolution of the planet more generally.  
43 Here, we develop the connection between the history of the fluid envelope and that of  
44 the silicate Earth. We use ideas about Earth accretion and insights that they yield for the  
45 origin and history of Earth’s volatiles. The focus is on Earth’s steam atmosphere, a  
46 unique reservoir in Earth history that links the energetic process of planetary accretion  
47 via giant impact to the emergence of Earth’s oceans and Hadean climate. We use “steam  
48 atmosphere” to describe any atmosphere prevented from condensation by internal heat  
49 in which  $\text{H}_2\text{O}$  is an important component, even those in which other gases (e.g.,  $\text{H}_2$ ) are  
50 more abundant by number.

51

52 A question closely related to the atmospheric composition of the early Earth is the oxygen  
53 fugacity ( $fO_2$ ) of the magma ocean from which the primordial atmosphere was outgassed  
54 (Elkins-Tanton, 2008). Plausible  $fO_2$  of magma ocean outgassing range from reducing  
55 ( $H_2$ -CO-rich) to oxidizing ( $H_2O$ - $CO_2$ -rich) atmospheres (Hirschmann, 2012). In analogy  
56 with the volatile inventory on modern Earth, prior work has assumed that the primordial  
57 atmosphere of early Earth was  $H_2O$ - $CO_2$ -rich (Abe and Matsui, 1988; Kasting, 1988).  
58 However, given a lack of knowledge about the oxygen fugacity of outgassing, reducing  
59 primordial outgassed atmospheres are difficult to rule out (Hirschmann, 2012). Models  
60 and measurements of early Earth oxygen fugacity yield contradictory evidence: whereas  
61 the  $fO_2$  of metal-silicate equilibration is necessarily reducing ( $\log fO_2 < IW-2$ ) due to the  
62 co-existence of metals (Wade and Wood, 2005) and implies  $H_2$ -CO-rich gas mixtures,  
63 the oldest terrestrial samples – Hadean zircons – are characterized by much higher  $fO_2$   
64 consistent with the modern (oxidized) terrestrial mantle, and suggest an  $H_2O$ - $CO_2$ -rich  
65 atmosphere (Trail et al., 2011). Some process apparently oxidized the silicate Earth  
66 during or shortly after core formation (Delano, 2001). Because the nature and timing of  
67 this process are as yet unclear, the chemical composition and oxidation state of Earth's  
68 early atmosphere remain unknown.

69

70 Here, by linking the hydrogen isotopic composition of the terrestrial oceans to the  
71 chemical composition of the primordial atmosphere, we articulate new constraints on  
72 the oxygen fugacity of magma ocean outgassing and subsequent processes on the Hadean  
73 Earth. Because this model is constrained by isotopic observations, it has the potential to  
74 yield new insights into early Earth evolution. Several features of early Earth evolution  
75 render the D/H composition of the oceans a reflection of the magma ocean and Hadean

76 atmospheric processes: (1) most of the water gained by Earth during planet formation  
77 was already accreted at the time of the Moon-forming giant impact and therefore  
78 participated in the terminal terrestrial magma ocean (Greenwood et al., 2018), (2) most  
79 (>75-98%) of the water initially dissolved in the magma ocean was outgassed into the  
80 atmosphere upon solidification, rendering the terrestrial water oceans the dominant  
81 exchangeable water reservoir on early Earth (Elkins-Tanton, 2008), and (3) the residence  
82 time of water in the terrestrial oceans with respect to the subduction and the deep water  
83 cycle is long, of order  $\sim 10^{10}$  years (van Keken et al., 2011). Together, these observations  
84 suggest that most hydrogen atoms currently residing in the oceans experienced the  
85 magma ocean and its aftermath and carry isotopic memory from these early epochs. The  
86 inference that the modern oceans reflect the isotopic composition of Hadean oceans –  
87 and, indeed, the bulk silicate Earth – is supported by measurements on Archean samples  
88 with D/H values for seawater identical to modern values to within a few percent (Pope  
89 et al., 2012). Percent-level variations in ocean D/H can arise due to exchange of water  
90 with the solid Earth due to plate tectonic processes (Kurokawa et al., 2018; Lécuyer et  
91 al., 1998) but here we are interested in large magnitude D/H variations ( $\sim 2x$ ) that can  
92 arise due to early atmospheric processes (see §4.3).

93

94 The isotopic composition of the oceans is determined by deuterium fractionation between  
95  $H_2O$  and  $H_2$  and the contrasting histories of these molecules in the planetary environment.  
96 Hydrogen ( $^1H$ ) and deuterium ( $^2H$ ) in water and molecular hydrogen experience distinct  
97 vibrational frequencies due to different bond strengths associated with the O-H and H-H  
98 stretch. These contrasting bonding environments result in deuterium being concentrated  
99 into water molecules while molecular hydrogen is deuterium-depleted at equilibrium,  
100 especially at low temperatures, i.e. isotopic fractionation occurs between the water ocean

101 and H<sub>2</sub>-rich atmosphere. Hence, even in the absence of HD/H<sub>2</sub> mass fractionation during  
102 the escape process (Zahnle et al., 1990), loss of molecular hydrogen from the early Earth  
103 can enrich planetary water in deuterium due to low-temperature equilibrium partitioning  
104 alone, i.e. light isotopes are concentrated into the escaping gas relative to the oceans.  
105 This equilibrium partitioning between molecules can result in isotopic evolution of the  
106 terrestrial oceans due to the contrasting fates of molecular hydrogen and water on early  
107 Earth. As on the modern Earth, water vapor is expected to be confined to the lower  
108 atmosphere via condensation at the tropopause (Wordsworth and Pierrehumbert, 2013b)  
109 and retained on the Hadean Earth while hydrogen in non-condensable forms (e.g., H<sub>2</sub>)  
110 can traverse the tropopause and undergo large-scale escape (Hunten, 1973). To the extent  
111 that the isotopic composition of escaping gas was distinct from that of oceans, isotopic  
112 evolution would have taken place.

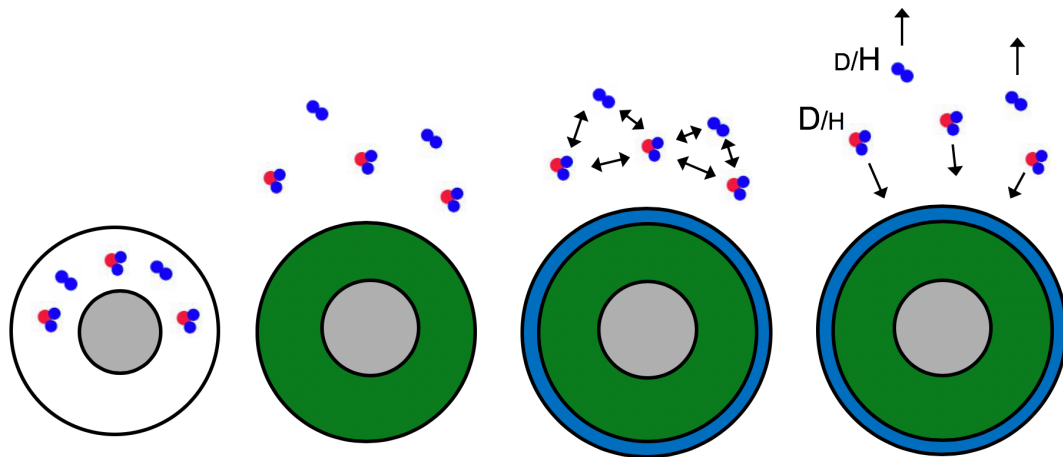
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114 That the hydrogen isotopic composition of the terrestrial oceans retains memory of early  
115 epochs permits its use in constraining early atmospheric fractionation processes, pending  
116 knowledge of the D/H of its source. Indeed, the source of the Earth's hydrosphere is  
117 commonly inferred from the D/H composition of the oceans in comparison with Solar  
118 System sources such as comets, asteroids (via meteorites), and the solar nebula (via  
119 Jupiter). Based on hydrogen and nitrogen isotopic evidence, the major terrestrial volatiles  
120 (C, N, H) are commonly inferred to be sourced primarily from carbonaceous chondrites  
121 (Alexander et al., 2012; Halliday, 2013; Marty, 2012). The close "match" (to within  
122 ~10%) between seawater D/H and the carbonaceous chondrite distribution peak – if not  
123 genetic – must be relegated to coincidence with a priori low probability (Lécuyer et al.,  
124 1998). Although a predominantly nebular source for terrestrial hydrogen with strong D/H  
125 fractionation has been explored (Genda and Ikoma, 2008), the nitrogen isotopic

126 evidence, in particular, argues against a predominantly nebular source for Earth's  
127 volatiles (Alexander et al., 2012; Marty, 2012). We adopt the common view that  
128 terrestrial seawater has a primarily carbonaceous chondritic source and show that the  
129 preservation of such a chondritic D/H signature in the terrestrial oceans places an upper  
130 limit on H<sub>2</sub> abundances on the Hadean Earth, requires oxidizing conditions for magma  
131 ocean outgassing (H<sub>2</sub>/H<sub>2</sub>O<0.3), and suggests a limited role for H<sub>2</sub> production via the  
132 iron-water reaction during terrestrial late accretion.

133

134 The outline of the paper is as follows. In §2, we estimate equilibration timescales between  
135 the magma ocean and primordial atmosphere and present calculations that relate the  
136 chemical composition of the outgassed atmosphere to the oxidation state of the magma  
137 ocean. In §3, we describe a climate model that considers greenhouse warming by the  
138 primordial H<sub>2</sub> inventory in the subsequent water ocean epoch and the timescales that  
139 govern its evolution. In §4, we describe the results of the model for the hydrogen isotopic  
140 evolution of the oceans due to equilibration and loss of a primordial H<sub>2</sub> inventory. In §5,  
141 we discuss the implications of these results for the oxidation state of the magma ocean  
142 and the oxidant for terrestrial late accretion, and in §6, we summarize and conclude. The  
143 envisioned sequence explored in this paper is summarized in Figure 1.



**Figure 1 – Behavior of hydrogen on Earth after the Moon-forming giant impact. (a)** After the giant impact, a deep magma ocean dissolves most of the hydrogen accreted to Earth, (b) crystallization of the magma ocean leads to primordial outgassing of most of exchangeable hydrogen with  $H_2/H_2O$  determined by oxygen fugacity of last equilibration with the magma ocean (§2.2), (c) condensation of the oceans and low-temperature ( $\sim 300$ - $500$ K)  $H_2O$ - $H_2$  D/H equilibration leads to deuterium-enrichments ( $H_2O$ ) and depletions ( $H_2$ ) in co-existing species, (d) retention of water via condensation coupled with loss to space of deuterium-depleted  $H_2$  via hydrodynamic escape leads to deuterium enrichment in the terrestrial oceans whose magnitude depends on the initial  $H_2/H_2O$  of the outgassed atmosphere.

144 **2. Magma ocean outgassing**

145 Magma oceans may arise in the early Solar System through various processes (Elkins-  
 146 Tanton, 2012). The Moon-forming giant impact extensively melts the silicate Earth and  
 147 leaves the accreting planet with  $\sim 99\%$  of its final mass (Pahlevan and Morbidelli, 2015).  
 148 During the ensuing magma ocean crystallization period, terrestrial water transitions from  
 149 being predominantly dissolved in the magma ocean to primarily outgassed into the steam  
 150 atmosphere, subsequent condensation of which forms the early terrestrial oceans (Elkins-



151 Tanton, 2008). This event is therefore considered the major volatile processing event in  
152 early Earth history, after which the abundance and initial distribution of terrestrial water  
153 is largely established and planetary processes transition from the accretionary to the  
154 geological regime. In this section, we first justify the use of equilibrium thermodynamics  
155 in calculating atmospheric compositions (§2.1) and then discuss the dependence of the  
156 outgassed gaseous composition on the redox state of the magma ocean at the time of last  
157 equilibration with the atmosphere (§2.2).

158

## 159 2.1. Magma ocean-atmosphere equilibration timescales

160 To justify the use of equilibrium thermodynamics, we first examine equilibration times  
161 between the magma ocean and the overlying atmosphere. On the modern Earth, the  
162 timescale for pCO<sub>2</sub> equilibration with the oceans is ~10<sup>2</sup> years (Archer et al., 2009), but  
163 no equivalent empirical estimate exists for magma ocean-atmosphere equilibration.  
164 Using a boundary-layer analysis (Hamano et al., 2013), we estimate the timescale for  
165 magma ocean-atmosphere equilibration assuming that diffusion through the magma  
166 surface boundary layer, rather than ascent of bubble plumes, dominates the equilibration  
167 process. We consider a schematic sequence in which thermal boundary layers form at  
168 the magma ocean surface and are peeled away by negative buoyancy. The equilibration  
169 timescale can then be estimated:

$$170 \quad \tau_{eq} = \tau_{BL} x N = (\delta_T^2 / \kappa) x (z / \delta_C) \quad (1)$$

171 where  $\tau_{BL}$  is the timescale for the formation of a thermal boundary layer, and  $N$  is the  
172 number of formation and buoyant destruction cycles before the entire magma ocean mass  
173 is processed through the surface boundary layer,  $\delta_T$  and  $\delta_C$  are the thermal and chemical  
174 boundary layer thicknesses, respectively,  $\kappa$  is the thermal diffusivity of the magma (cm<sup>2</sup>

175  $s^{-1}$ ), and  $z$  is magma ocean depth. We can express the thickness of the chemical boundary  
176 layer in terms of the thickness of the thermal boundary layer via scaling:  $(\delta_C/\delta_T) =$   
177  $(D/\kappa)^{1/2}$  where we take  $D$  as the atomic diffusivity for water in magma ( $cm^2 s^{-1}$ ). For  
178 parameter choices, we adopt a thermal diffusivity ( $\kappa=2 \times 10^{-3} cm^2 s^{-1}$ ) appropriate to  
179 peridotite liquid (Leshner and Spera, 2015) and an atomic diffusivity for water in magma  
180 ( $D_w=1.4 \times 10^{-4} cm^2 s^{-1}$ ) appropriate for a basaltic magma with  $\sim 1$  wt% water at 2000 K  
181 (Zhang et al., 2007). Substitution of these parameters yields an estimate for equilibration  
182 timescales:

$$183 \quad \tau_{eq} = 3 \times 10^3 \text{ years} \times \left( \frac{\delta_T}{1 \text{ cm}} \right) \times \left( \frac{z}{500 \text{ km}} \right) \quad (2)$$

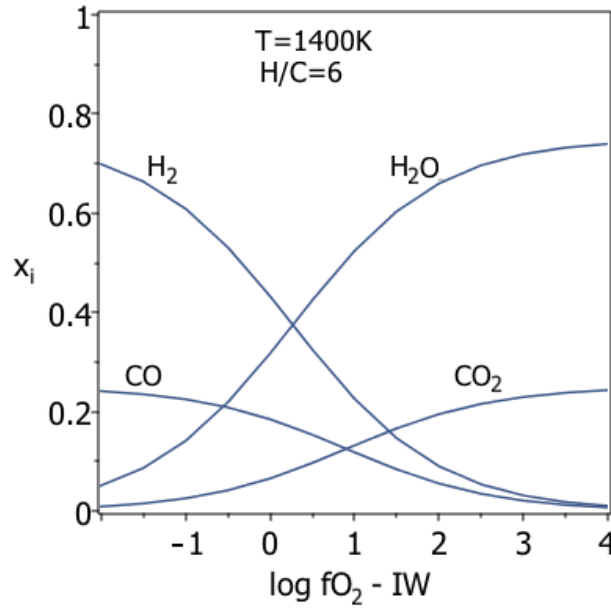
184 This result suggests that equilibration with the atmosphere by processing the magma  
185 ocean through a chemical boundary layer is rapid relative to cooling timescales, which  
186 are estimated at  $\sim 10^6$  years (Lebrun et al., 2013). Although more work is required to  
187 quantify the competing role of magma outgassing via bubble plumes in accommodating  
188 the super-saturation near the surface (Ikoma et al., 2018), this treatment suggests that  
189 equilibration via boundary layer diffusion alone may be sufficiently rapid to motivate the  
190 equilibrium assumption.

191

## 192 2.2. Primordial atmospheric compositions

193 A fundamental parameter governing equilibrium compositions of outgassed atmospheres  
194 is the oxygen fugacity ( $fO_2$ ). As long as the redox buffering capacity of the magma ocean  
195 exceeds that of the outgassed atmosphere, the magma ocean determines the  $fO_2$  of the  
196 atmosphere with which it equilibrates (Hirschmann, 2012). Volatile outgassing from the  
197 magma ocean continues until the end of magma ocean crystallization, when the  
198 formation of a meters-thick solid chill crust effectively isolates the newly formed

199 atmosphere from exchange with silicate Earth. Therefore, the initial composition of the  
200 atmosphere during the Hadean is dictated by last equilibration with the magma ocean.  
201  
202 Despite its importance to planetary evolution, the  $fO_2$  at the magma ocean-atmosphere  
203 interface is not well-constrained. There are two end-members: (1) in analogy with the  
204 modern Earth mantle, the magma ocean-steam atmosphere system may be chemically  
205 oxidized ( $\log fO_2 \sim QFM$ ), with water vapor and carbon dioxide dominant (Abe and  
206 Matsui, 1988; Kasting, 1988). However, (2) the terrestrial magma ocean – having held  
207 metallic droplets in suspension – may also be much more reducing ( $\log fO_2 \sim IW-2$ ) at the  
208 surface where equilibration with the atmosphere takes place. Under such reducing  
209 conditions,  $H_2$  and  $CO$  are the dominant H- and C-bearing gaseous species (Fig. 2). A  
210 remarkable feature of magma oceans is that the expected range of possible  $fO_2$  values  
211 spans the transition from the reducing ( $H_2$ - $CO$ -rich) to oxidizing ( $H_2O$ - $CO_2$ -rich) gas  
212 mixtures, indicating a rich volatile-processing history during planetary accretion.



**Figure 2 – High-temperature equilibrium outgassed atmospheres.**

The mole fraction of vapor species is calculated as a function of oxygen fugacity ( $fO_2$ ) relative to the iron-wüstite (IW) buffer at an equilibrium temperature of 1,400 K. Parameters for the IW buffer are given in (Frost, 1991). Thermodynamic data for gaseous species ( $H_2O$ - $H_2$ - $CO$ - $CO_2$ ) are adopted from (Chase et al., 1985). The lower end of the range of  $\log fO_2$  (IW-2) characterizes oxygen fugacity of a magma ocean in equilibrium with a suspension of metallic droplets, whereas the upper end of the range (IW+4 $\approx$ QFM) corresponds to the redox state of oxidized modern Earth mantle.

213 Calculation of climate (§3) requires specification of the mass and molecular composition  
 214 of the atmosphere. Whereas an oxidizing ( $H_2O$ - $CO_2$ -rich) outgassed gaseous envelope  
 215 experiences minimal chemistry upon cooling and maintains its molecular composition,  
 216 evolution of a reducing ( $H_2$ - $CO$ -rich) gaseous envelope may involve more significant  
 217 transformations. As the envelope cools following magma ocean crystallization, the

218 reaction ( $3\text{H}_2+\text{CO}\rightleftharpoons\text{CH}_4+\text{H}_2\text{O}$ ) shifts to the right, potentially converting the outgassed  
219 gas mixture into a methane-rich atmosphere before quenching and condensation of the  
220 oceans (Schaefer and Fegley, 2010). Nevertheless, we consider  $\text{CH}_4$  to be – at most – a  
221 transient molecule in the primordial atmosphere because it is unstable with respect to UV  
222 photolysis via Lyman  $\alpha$  emission. The photolysis of methane under the influence of the  
223 UV flux of the young Sun occurs rapidly, at a rate of  $\sim 1$  bar/Myrs (Kasting, 2014), with  
224 the carbon oxidized to  $\text{CO}/\text{CO}_2$  and the hydrogen reverting to molecular form before  
225 escaping, analogous to the modern atmosphere of Titan. The photochemical stability of  
226  $\text{H}_2$  suggests that the reducing power inherited from the magma ocean is primarily carried  
227 by, and lost via the escape of, molecular hydrogen.

228

### 229 **3. Climate model**

230 In this section, we describe a climate model used to calculate surface temperatures due  
231 to  $\text{H}_2$  greenhouse warming, which we use in the next section (§4) to determine the relative  
232 behavior of hydrogen and deuterium on the Hadean Earth. Following magma ocean  
233 crystallization,  $\text{H}_2\text{O}$  in the outgassed atmosphere becomes the early oceans. The partial  
234 pressure of  $\text{H}_2\text{O}$  is then determined by vapor pressure equilibrium at surface temperatures  
235 determined by solar radiative balance. Hence, soon after magma ocean crystallization,  
236 the Hadean Earth will relax into a solar-powered climate with early water oceans co-  
237 existing with a massive non-condensable primordial atmosphere. While oxidizing ( $\text{H}_2\text{O}$ -  
238  $\text{CO}_2$ -rich) ocean climates have previously been described (Wordsworth and  
239 Pierrehumbert, 2013b), equivalent reducing ( $\text{H}_2\text{O}$ - $\text{H}_2$ -rich) ocean climates have not. To  
240 describe the evolution of the earliest Hadean climate, we adopt a 2-component ( $\text{H}_2\text{O}$ - $\text{H}_2$ )  
241 chemical model for the atmosphere and ocean and take the oxygen fugacity of last  
242 equilibration between the magma ocean and primordial atmosphere (and hence outgassed

243 H<sub>2</sub>/H<sub>2</sub>O) as a free parameter. Because water condenses in the lower atmosphere and is  
244 retained but molecular hydrogen can escape, the free parameter governing early climate  
245 can be expressed as the inventory of molecular hydrogen in the atmosphere (pH<sub>2</sub>).

246

247 Due to water condensation and cold-trapping in the lower atmosphere and collision-  
248 induced infrared opacity of H<sub>2</sub> at high (>0.1 bar) pressures, the emission level in H<sub>2</sub>O-  
249 H<sub>2</sub> model atmospheres is determined by the opacity of H<sub>2</sub> (Wordsworth, 2012). The  
250 emission temperature (T<sub>E</sub>) is given by top-of-the-atmosphere radiative balance with the  
251 early Sun:

$$252 \quad \frac{L}{4}(1 - A) = \sigma_{SB}T_E^4 \quad (3)$$

253 with L the solar constant, A the visible bond albedo, and  $\sigma_{SB}$  the Stefan-Boltzmann  
254 constant. For L=10<sup>3</sup> W/m<sup>2</sup> appropriate for the early Sun and a bond albedo A=0.3, an  
255 emission temperature of 235 K is obtained for the primordial Earth. For simplicity, solar  
256 luminosity in these calculations is held constant. Results are qualitatively similar for a  
257 range of emission temperatures (215-255 K), as might be expected based on an evolving  
258 Sun, cloud feedback and/or by adjusting the planetary albedo to account for increased  
259 Rayleigh scattering in thicker atmospheres. These effects are known to alter the radiation  
260 budget by tens of percent (Gough, 1981; Wordsworth, 2012). At infrared wavelengths,  
261 the optical depth unity surface of a pure H<sub>2</sub> atmosphere has been calculated for a several  
262 Earth-mass planet (g=20 m/s<sup>2</sup>, T<sub>ph</sub>=100 K) and is ~0.2 bars (Wordsworth, 2012).  
263 Combining the expressions for photospheric pressure (P<sub>ph</sub>∝g/κ) and collision-induced  
264 opacity (κ∝ρ) of an ideal gas (ρ∝P<sub>ph</sub>/T<sub>ph</sub>) yields a scaling relation for photospheric  
265 pressure, gravitational acceleration, and photospheric temperature (P<sub>ph</sub>∝g<sup>1/2</sup>T<sub>ph</sub><sup>1/2</sup>).  
266 Applying this scaling relation to the emission temperatures (T<sub>ph</sub>=235K) relevant to the  
267 Hadean Earth (g=9.8m/s<sup>2</sup>) yields an H<sub>2</sub> photospheric pressure ~0.21 bars, which we

268 adopt. Since the atmosphere at the emission level is cold and dry, this pressure –  
269 appropriate for a pure H<sub>2</sub> atmosphere – is taken as the emission pressure to which a moist  
270 adiabatic structure (see below) must be stitched.

271

272 Planetary surface temperature is controlled by the structure of the troposphere. At the  
273 base of the troposphere, vapor pressure equilibrium with the ocean controls the water  
274 vapor abundance. We assume a troposphere saturated in water vapor throughout.  
275 Accordingly, the partial pressure of water vapor is given by  $p_{H_2O} = \exp(-\Delta G/RT)$ ,  
276 with  $\Delta G = \Delta H - T\Delta S$  and  $\Delta H = 40.58$  kJ/mol,  $\Delta S = 0.1082$  kJ/mol.K (Chase et al., 1985). As  
277 in any multi-component atmosphere, the inventory of one gas influences the partial  
278 pressure of other gases through vertical redistribution (Wordsworth and Pierrehumbert,  
279 2013a). The total pressure is assumed given by the expression  $p_T = \sigma_{H_2O} g \mu / \mu_{H_2O} x_{H_2O}$   
280 with  $\sigma_{H_2O}$  the surface density of water vapor,  $g$  the gravitational acceleration at the  
281 surface,  $\mu$  the mean molecular weight, and  $\mu_{H_2O}$  and  $x_{H_2O}$  the atomic mass and the mole  
282 fraction of water vapor. This expression is strictly valid in a well-mixed atmosphere,  
283 which we take as an adequate approximation for an H<sub>2</sub>O-H<sub>2</sub> atmosphere. With the above  
284 relations, we write an expression for the entropy of atmospheric gas in contact with the  
285 oceans, which we use to calculate surface temperature. The entropy of the troposphere is  
286 that of the basal gaseous mixture:

$$287 \quad S(T) = \sum_i x_i s_i(T) - R \sum_i x_i \ln x_i - R \ln P \quad (4)$$

288 which is the expression for entropy of an ideal mixture of ideal gases, with the first term  
289 the sum over species as pure gases at standard pressure, the second term an entropy of  
290 mixing term deriving from the fact that the gas parcel is a mixture of randomly distributed  
291 gas molecules, and the third term a pressure correction due to the volume available to  
292 each molecule. In this way, the entropy of the atmospheric parcel at any temperature (T),

293 pressure (P) and composition ( $x_i$ ) can be calculated. We consider ideal gas theory as an  
294 adequate approximation for the primordial atmosphere because intermolecular distances  
295 are large relative to the size of the molecules. The thermodynamic data for the entropy  
296 of pure substances is also taken from standard sources (Chase et al., 1985).

297

298 Since a moist adiabat is also isentropic, the tropospheric entropy can be used to relate the  
299 conditions at the base to those characterizing the radiative emission level where the mode  
300 of energy transport transitions to radiation. A major influence on tropospheric structure  
301 is the condensation of water vapor into clouds through adiabatic expansion and cooling.  
302 The specific entropy at the radiative emission level at the top of the troposphere must  
303 therefore consider condensates:

$$304 \quad S(T) = F_v S_v(T) + F_L S_L(T) \quad (5)$$

305 with  $F_v$  and  $F_L$  the fraction of total molecules in the parcel in the gas and condensates,  
306 respectively. Atmospheric P-T paths determined by calculating pseudoadiabats (i.e., with  
307 rainout) are similar to adiabatic equivalents (i.e. with suspended condensates) (Ingersoll,  
308 1969), permitting use of a two-phase isentrope even when rainout might be expected and  
309 upper tropospheric opacity determined by gas-phase ( $H_2$ ) opacity alone. The procedure  
310 for calculating surface temperature then entails: (1) an initial estimate for the surface  
311 temperature, yielding the partial vapor pressure of water vapor and, with a given  $H_2$   
312 inventory, the total atmospheric surface pressure, (2) calculation of the entropy of the  
313 convective atmosphere, (3) evaluation of thermodynamic state, including temperature,  
314 of the atmospheric parcel at a pressure of 0.21 bars representing the radiative emission  
315 level. In this way, each value of surface temperature corresponds to a value of emission  
316 temperature, and iteration allows identification of the surface temperature corresponding  
317 to the emission temperature required by top-of-the-atmosphere radiative balance



318 (Equation 3). Because the convective troposphere is isentropic, we can solve for the  
319 surface temperature as a function of emission temperature without explicitly resolving  
320 the vertical structure in the intervening atmosphere. In this way, we calculate surface  
321 temperatures in the all-troposphere approximation (Pierrehumbert, 2010) and iterate to  
322 find solutions.

323

## 324 **4. Results**

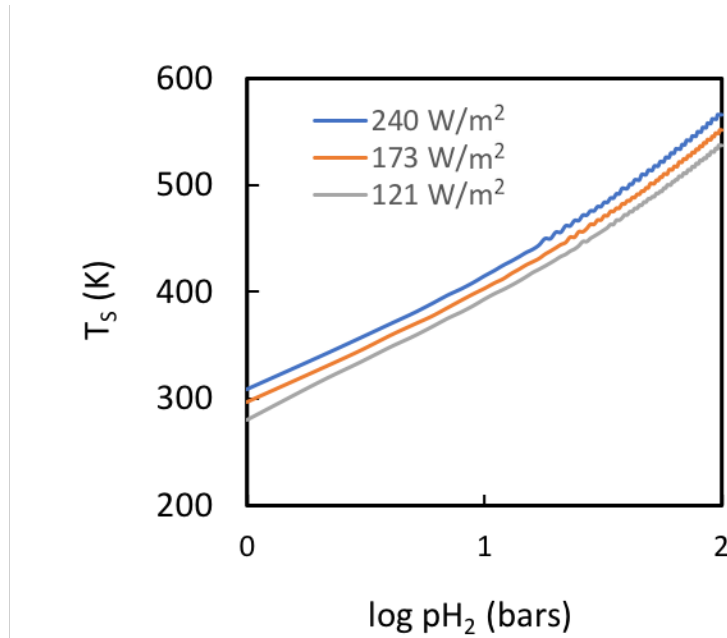
325 In this section, we describe the results of the climate model (§4.1) and use the results to  
326 compare the timescales for ocean-atmosphere isotopic equilibration with atmospheric  
327 escape (§4.2) to motivate equilibrium isotopic partitioning to describe the behavior of  
328 hydrogen and deuterium on the early Earth. Finally, we discuss the hydrogen isotopic  
329 evolution of the Hadean oceans and the derived upper limit on  $p\text{H}_2$  on the early Earth  
330 (§4.3).

331

### 332 4.1. Equilibrium average surface temperatures

333 Greenhouse warming by  $\sim 1\text{-}10^2$  bars of equivalent  $\text{H}_2$  on the Hadean Earth yields surface  
334 temperatures  $\sim 300\text{-}550$  K (Fig. 3). Despite some uncertainty arising from cloud feedback  
335 or higher albedo due to greater Rayleigh scattering, multi-bar atmospheric  $\text{H}_2$  inventories  
336 are sufficient to keep Earth out of a snowball state via  $\text{H}_2$  opacity alone. The calculation  
337 assumes radiative transfer in the terrestrial planet regime, where a sufficient fraction of  
338 the visible flux penetrates the atmosphere and powers convection throughout the  
339 troposphere, whereas for  $p\text{H}_2 > 20$  bars visible photons might not penetrate, creating a  
340 deep atmospheric structure governed by geothermal rather than solar heating as in  
341 interior of the giant planets (Wordsworth, 2012). Calculated surface temperatures for  
342  $p\text{H}_2 > 20$  bars are therefore upper limits. These calculations suggest that the dominant

343 control on the surface temperatures in the Hadean is the atmospheric H<sub>2</sub> inventory and  
344 that the earliest climate on Earth was warm and governed by the physics of atmospheric  
345 escape.



**Figure 3 – Ocean surface temperature as a function of primordial H<sub>2</sub> inventory.** Temperatures are shown for different values of the outgoing longwave radiation (OLR) flux for emission temperatures of 215, 235, and 255 K, corresponding to a range (0-0.5) of planetary bond albedos. The primordial planetary climate depends primarily on the H<sub>2</sub> inventory. The dominant control on earliest climate in these scenarios is the physics of atmospheric escape.

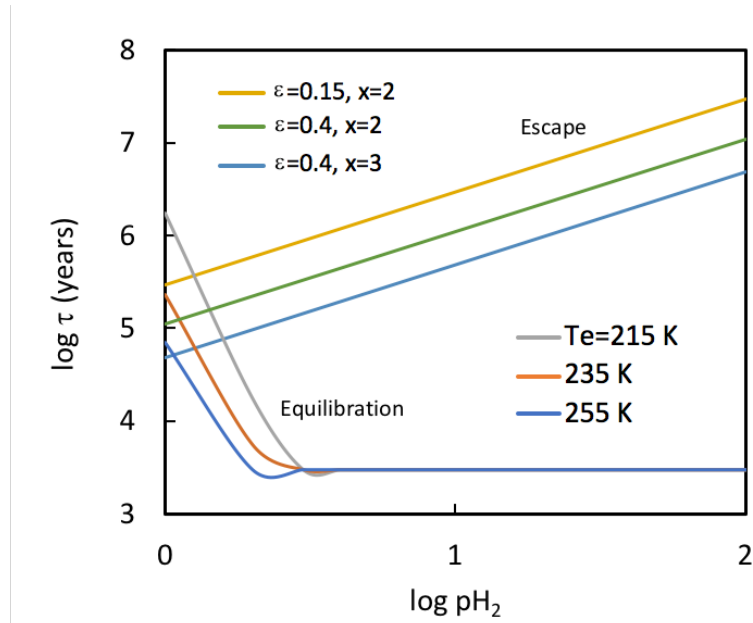
346 4.2. Atmospheric equilibration and escape timescales

347 Surface temperatures derived from the climate model can be used to assess H isotopic  
348 equilibration between the oceans and the atmosphere. A key comparison is between the  
349 equilibration time and the residence time of atmospheric H<sub>2</sub>. H<sub>2</sub>O-H<sub>2</sub> isotopic exchange  
350 may be rate-limited by exchange in the atmosphere, because H<sub>2</sub> dissolves negligibly in

351 the water oceans. Isotope exchange between water vapor and molecular hydrogen is rapid  
352 and highly temperature-dependent (Lécluse and Robert, 1994). After the condensation  
353 of the steam into water oceans (§3), a significant H<sub>2</sub> inventory (>10-100 bars H<sub>2</sub>) and  
354 high surface temperatures (>400-550 K, Fig. 3) result in extremely rapid atmospheric  
355 reactions such that the timescale for ocean-atmosphere equilibration is limited by the rate  
356 of ocean evaporation and circulation of water through the atmosphere in a hydrological  
357 cycle. This timescale ( $\equiv$  depth of oceans/evaporation rate) is rapid ( $\approx 10^3$  years), and  
358 nominally independent of H<sub>2</sub> inventory and surface temperature (e.g. p<sub>H<sub>2</sub></sub>>3 bars, Fig.  
359 4). Ocean-atmosphere isotopic equilibration during the Hadean is therefore expected to  
360 initially proceed rapidly. As the primordial climate evolves due to H<sub>2</sub> escape, reaction  
361 rates decline until ocean-atmosphere equilibration becomes rate-limited by exchange  
362 reactions in the atmosphere. In this regime, the equilibration time between the oceans  
363 and atmosphere is given by:

$$364 \quad \tau_{AO} = (\tau_{ex}/\tau_{res}) \times \tau_{cir} \quad (6)$$

365 with  $\tau_{ex}$  the isotopic exchange time between atmospheric H<sub>2</sub> and H<sub>2</sub>O,  $\tau_{res}$  the residence  
366 time of atmospheric water vapor, and  $\tau_{cir}$  the timescale to circulate the entire oceans  
367 through the atmosphere via evaporation and precipitation in a global hydrological cycle  
368 (Genda and Ikoma, 2008). As the atmospheric H<sub>2</sub> inventory is lost, the ocean-atmosphere  
369 equilibration time evolves for two reasons: (1) the isotopic exchange timescale ( $\tau_{ex}$ )  
370 increases due to the cooling temperatures, (2) the residence time of water vapor in the  
371 atmosphere ( $\equiv$  depth of equivalent water layer/precipitation rate) decreases due to the  
372 lower water vapor pressures associated with lower temperature. Both effects prolong the  
373 equilibration time between the ocean and atmosphere, which increases rapidly at low  
374 temperatures (p<sub>H<sub>2</sub></sub><3 bars, Fig. 4), approaching  $\sim 1$  Myrs for H<sub>2</sub> inventories of  $\sim 1$  bar.



**Figure 4 – Ocean-atmosphere isotopic equilibration and hydrodynamic escape versus primordial H<sub>2</sub> inventories.** Isotopic equilibration between the water oceans and a primordial H<sub>2</sub>-rich atmosphere is calculated for three values of the top-of-the-atmosphere emission temperature ( $T_e$ ), as a function of total atmospheric H<sub>2</sub> inventory. At high H<sub>2</sub> inventories (e.g. 10-100 bars equivalent H<sub>2</sub>), the equilibrium temperatures are high ( $\sim$ 400-500K, Fig. 3), and isotopic exchange reactions are extremely fast. Parameters characterizing the escape times are the thermal efficiency ( $\epsilon$ ) and planetary EUV absorption radii ( $x$ ) (§4.2). Timescales for ocean-atmosphere equilibration are generally shorter than the residence time of atmospheric H<sub>2</sub> with respect to escape suggesting continuous equilibration during the escape process.

375 Ocean-atmosphere isotopic equilibration requires that the equilibration time be shorter  
 376 than the residence time of H<sub>2</sub> in the atmosphere. To determine whether or not this is the  
 377 case, we calculate extreme ultraviolet (EUV) powered escape rates (Watson et al., 1981)  
 378 for initial H<sub>2</sub> inventories, and adopt the assumption that loss to space is the sole H<sub>2</sub> sink,

379 as expected on the prebiotic Earth. Escape rates can be calculated using the energy-  
380 limited approximation:

$$381 \quad \phi_{H_2} = \frac{\epsilon F_{EUV}}{4R_p} x^2 \quad (\text{Pa s}^{-1}) \quad (7)$$

382 with  $\phi$  the hydrogen escape flux expressed as atmospheric pressure loss per unit time,  
383  $F_{EUV}$  the extreme ultraviolet flux of the young sun,  $\epsilon$  the thermal efficiency or the fraction  
384 of incident EUV used to power the planetary wind ( $\epsilon=0.15-0.4$ ),  $R_p$  the planetary radius,  
385 and  $x(\equiv R_{EUV}/R_p)$  the effective EUV absorption radius in planetary radii ( $x=2-3$ ), which  
386 characterizes the distended nature of EUV absorption in escaping atmospheres.  $F_{EUV}$  is  
387 assumed equal to 100x the modern extreme ultraviolet flux, i.e.  $464 \text{ erg cm}^{-2} \text{ s}^{-1}$  (Ribas  
388 et al., 2005). Other parameter choices ( $\epsilon, x$ ) are adopted from hydrodynamic calculations  
389 of hydrogen-rich atmospheres exposed to comparable EUV fluxes (Erkaev et al., 2016).  
390 Calculated  $H_2$  residence times ( $=p_{H_2}/\phi_{H_2}$ ) are  $\approx 10^6-10^7$  years for atmospheric inventories  
391  $\approx 10-100$  bars (Fig. 4). Timescales for ocean-atmosphere equilibration are generally  
392 shorter than residence times of  $H_2$  in the atmosphere, suggesting continuous equilibration  
393 with ocean-atmosphere quenching occurring at  $H_2$  inventories of only  $\sim$ a few bars.

394

#### 395 4.3. Hydrogen isotopic evolution of the Hadean oceans

396 Finally, we quantify the behavior of hydrogen ( $^1H$ ) and deuterium ( $^2H$ ) as tracers of early  
397 Earth evolution to articulate constraints on the chemical composition of the primordial  
398 atmosphere. In brief, water vapor is retained via condensation but hydrogen in non-  
399 condensable form (e.g.  $H_2$ ) interacts with water vapor in the lower atmosphere and is  
400 transported to the upper atmosphere and lost to space. Equilibrium D/H partitioning  
401 between water and hydrogen is calculated from standard prescriptions (Richet et al.,

402 1977).<sup>1</sup> While this prescription strictly relates water vapor to molecular hydrogen, it can  
403 also be used to characterize equilibrium between water oceans and H<sub>2</sub>-rich atmospheres  
404 because the vapor pressure isotope effect relating liquid water to water vapor is an order  
405 of magnitude smaller and can therefore be neglected to first order. H<sub>2</sub>O-H<sub>2</sub> equilibration  
406 is among the largest equilibrium fractionations between two molecules in nature, with a  
407 clearly resolvable magnitude at planetary temperatures. D/H composition of planetary  
408 oceans therefore reflects the mass of early H<sub>2</sub> reservoirs and the temperatures of isotopic  
409 equilibration. To the extent that atmospheric H<sub>2</sub> reservoirs comparable to the terrestrial  
410 oceans were present, the D/H composition of the hydrosphere could have evolved  
411 dramatically due to equilibrium partitioning and removal of the isotopically light H<sub>2</sub>.

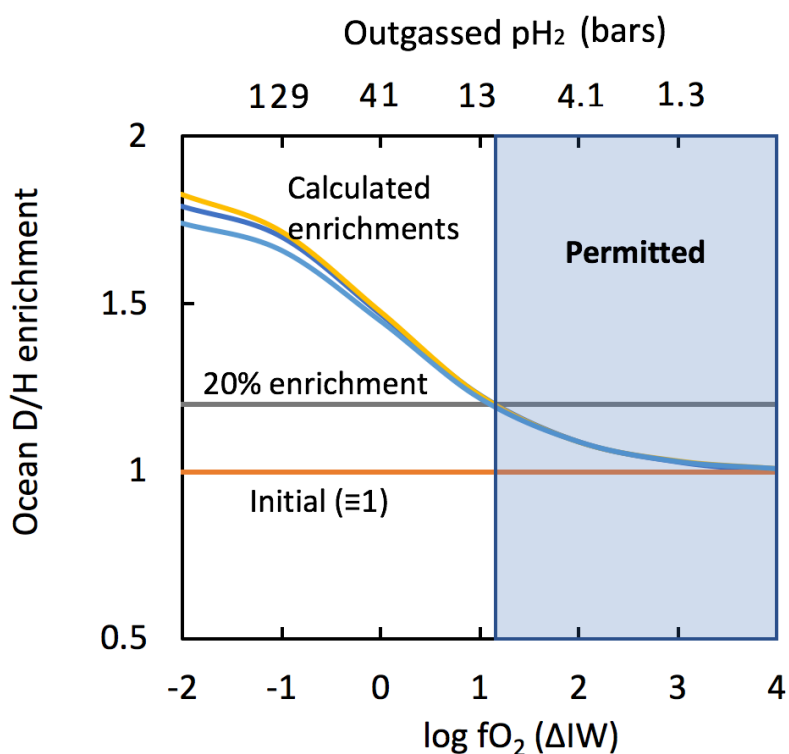
412

413 The hydrogen isotopic evolution of Hadean oceans can be calculated using equilibrium  
414 ocean-atmosphere partitioning. Following magma ocean crystallization, most of the  
415 water vapor condenses into oceans, while most H<sub>2</sub> partitions into the atmosphere,  
416 generating the earliest climate via radiative balance with the Sun (§3). The magnitude of  
417 greenhouse warming for a freshly outgassed atmosphere is significant: only a few percent  
418 of the outgassed volatile inventory need be in the form of molecular hydrogen to prevent  
419 a snowball Earth and to stabilize a water ocean climate via H<sub>2</sub> greenhouse warming alone  
420 (Fig. 3). The existence of such an early greenhouse climate permits isotopic equilibration  
421 between the ocean and atmosphere (§4.2) with the temperature-dependent partitioning  
422 between reservoirs determined self-consistently via climate with a given H<sub>2</sub> inventory  
423 (§3). As the atmospheric H<sub>2</sub> inventory is depleted via escape, the greenhouse effect also  
424 diminishes, accentuating the temperature-dependent partitioning between ocean and  
425 atmosphere. Hence, deuterium is further concentrated into the oceans due to the cooling

---

<sup>1</sup> The exchange reaction is  $\text{H}_2\text{O} + \text{HD} \rightleftharpoons \text{HDO} + \text{H}_2$  with equilibrium constant  $K(T) = 1 + 0.22 \cdot (10^3/T)^2$ .

426 radiative balance accompanying H<sub>2</sub> loss. In this way, D/H evolution of the oceans can be  
427 calculated as an equilibrium distillation sequence, converging to a value determined by  
428 the initial inventory of atmospheric H<sub>2</sub> (Fig. 5). Water molecules have a strong tendency  
429 to concentrate deuterium, and, to the extent that hydrogen in other forms (e.g. H<sub>2</sub>) are  
430 significant species in the primordial atmosphere, highly pronounced D/H enrichments  
431 (~2x) can arise from processing in this epoch, an enrichment not evident in the isotopic  
432 record of Earth (see §5). Such enrichment can be minimized for a very oxidizing (H<sub>2</sub>-  
433 poor) atmospheres outgassed from the magma ocean. The D/H composition of the water  
434 oceans immediately following primordial atmosphere evolution thereby constrains the  
435 H<sub>2</sub>/H<sub>2</sub>O ratio of Earth's outgassed atmosphere to <0.3 and logfO<sub>2</sub> of last equilibration to  
436 >IW+1. For a given H<sub>2</sub>O abundance, e.g., 1-2 modern ocean equivalents (Korenaga,  
437 2008), upper limits on the initial atmospheric H<sub>2</sub> inventory of Earth (pH<sub>2</sub><10-20 bars)  
438 can be derived (Fig. 5).



**Figure 5 – Ocean deuterium enrichment versus oxygen fugacity of primordial outgassing.** Oxygen fugacity determines the  $H_2/H_2O$  of the outgassed atmosphere. Oxidizing conditions ( $\log fO_2 \sim IW+4 \approx QFM$ ) lead to nearly “pure” steam atmospheres ( $H_2/H_2O < 0.01$ , see Fig. 2) and minimal D-enrichment in the oceans. More reducing conditions for outgassing (e.g.  $\log fO_2 < IW$ ) generate higher  $H_2/H_2O$  values ( $> 1$ ) and stronger ocean deuterium enrichments ( $> 1.5-2x$ ) due to equilibrium partitioning and loss of large quantities of isotopically light  $H_2$ . Minimal ( $< 20\%$ ) D-enrichment in terrestrial water relative to the source (§5) constrains the  $H_2/H_2O$  of the primordial atmosphere to  $< 0.3$  and therefore  $\log fO_2$  of outgassing to  $> IW+1$  (Permitted region).  $H_2$  abundances can be expressed in absolute terms (top axis) by fixing water abundances to the terrestrial ocean reservoir ( $pH_2O = 270$  bars). Temperatures of equilibration are calculated via the climate model (§3), with enrichment curves representing different emission temperatures ( $T_e = 215, 235, 255K$ ) demonstrating the robustness of the result.



439 Here, we have only considered ocean-atmosphere partitioning followed by the escape of  
440 atmospheric H<sub>2</sub> with no isotopic fractionation during the escape process. Consideration  
441 of kinetic processes (e.g. HD/H<sub>2</sub> mass-fractionation) will fractionate H isotopes further  
442 (Zahnle et al., 1990) and cause additional oceanic deuterium enrichment for a given H<sub>2</sub>  
443 inventory (Genda and Ikoma, 2008). The constraints on primordial H<sub>2</sub> abundances shown  
444 in Figure 5 are, therefore, conservative upper limits.

445

## 446 **5. Discussion**

447 The molecular composition of the Earth's primordial atmosphere is not well-constrained.  
448 Nevertheless, on the basis of their isotopic compositions, the Earth's major volatiles (H,  
449 N, C) are thought to be sourced primarily from the carbonaceous chondrites (Alexander  
450 et al., 2012; Halliday, 2013; Marty, 2012). This widely-held view of the source of major  
451 terrestrial volatiles requires preservation of the source signature in the terrestrial oceans  
452 and implies minimal D-enrichment via equilibration and escape of primordial H<sub>2</sub>. To  
453 quantify the constraint that this comparison places on primordial outgassing, we compare  
454 the isotopic composition ( $\delta D = -25 \pm 5\%$ ) of the Archean oceans (Pope et al., 2012) with  
455 the lowest bulk chondritic values ( $\delta D = -226 \pm 4\%$ ) measured to date (Alexander et al.,  
456 2012). On this basis, a primarily chondritic source for terrestrial water requires minimal  
457 (<20%) deuterium-enrichment via H<sub>2</sub> loss. According to our model calculations, this  
458 level of isotopic preservation requires most outgassed hydrogen from the terrestrial  
459 magma ocean to appear in the form of water vapor (H<sub>2</sub>/H<sub>2</sub>O < 0.3, Fig. 5) and requires the  
460 last equilibration with the magma ocean to be characterized by logfO<sub>2</sub> higher than IW+1.  
461 By connecting the conditions of outgassing to the observable isotopic signatures of  
462 ancient seawater, we can articulate new constraints on the composition of the Earth's

463 primordial atmosphere. These results show that reducing gases such as H<sub>2</sub> and CH<sub>4</sub> made  
464 up only a minor fraction of the Earth's outgassed atmosphere and require that the late-  
465 stage magma ocean was already oxidized by the time of last equilibration with the  
466 atmosphere. In this section, we discuss the implications of these results for redox  
467 evolution of the magma ocean (§5.1) and the oxidant involved in terrestrial late accretion  
468 (§5.2).

469

#### 470 5.1. The redox state of the terrestrial magma ocean

471 The redox state of a magma ocean determines both the chemical composition of the  
472 outgassed atmosphere and the isotopic composition of the water ocean following  
473 primordial H<sub>2</sub> loss. Oxidizing magma oceans outgas water-rich primordial atmospheres,  
474 which condense into oceans, experiencing minimal hydrogen escape and deuterium  
475 enrichment. Reducing magma oceans, by contrast, outgas substantial quantities of  
476 hydrogen in non-condensable form (e.g. H<sub>2</sub>) in addition to water molecules whose  
477 equilibration with the early water oceans can strongly (~2x) enrich water oceans in D/H,  
478 a feature not evident in the isotopic record of Earth (Pope et al., 2012). The persistence  
479 of a “chondritic” signature in the terrestrial oceans requires a low atmospheric H<sub>2</sub>/H<sub>2</sub>O  
480 at the time of last equilibration with a magma ocean, a ratio that acts as an oxybarometer,  
481 implying relatively oxidizing conditions ( $\log f_{\text{O}_2} > \text{IW} + 1$ ) for primordial outgassing.  
482 Given that the vigorously convective magma ocean initially held metallic droplets in  
483 suspension (Stevenson, 1990) and was therefore more chemically reducing ( $\log f_{\text{O}_2} < \text{IW} -$   
484 2) at early times, these results suggest that the silicate Earth was oxidized *during* the  
485 evolution of the magma ocean. Three mechanisms have been discussed for this  
486 primordial oxidation: (1) the terrestrial magma ocean could have been oxidized via Fe  
487 disproportionation at high pressure ( $3\text{Fe}^{+2} \rightarrow 2\text{Fe}^{+3} + \text{Fe}^0$ ) with separation of the newly

488 generated metallic iron to the core leaving an oxidized mantle residue (Wade and Wood,  
489 2005), (2) the primordial atmosphere was reducing (H<sub>2</sub>-rich) but the process of H<sub>2</sub> escape  
490 during the lifetime of the magma ocean oxidized both silicate Earth and co-existing  
491 atmosphere (Hamano et al., 2013), and (3) the (Fe<sup>+3</sup>/Fe<sup>+2</sup>) value of the terrestrial magma  
492 ocean was low (~0.01) but the more incompatible nature of ferric iron (Fe<sup>+3</sup>) in mantle  
493 minerals enriched it in evolving liquids such that the late-stage magma was more  
494 oxidizing than the magma ocean at the outset of crystallization (Schaefer et al. in prep).  
495 The relative importance of these three processes for the redox evolution of magma oceans  
496 is subject to future study. For now, we conclude that the oxidation of the silicate Earth  
497 occurred during the crystallization of the magma ocean, independently corroborating the  
498 conclusion from geological data for early (>4 Gya) establishment of oxidizing conditions  
499 in the silicate Earth (Trail et al., 2011).

500

## 501 5.2. The oxidant for terrestrial late accretion

502 Before outgassing of the primordial atmosphere, the magma ocean potentially facilitates  
503 the last major episode of core formation via separation of metallic droplets accompanying  
504 deep magma ocean convection on rapid (~10<sup>2</sup> year) timescales (Stevenson, 1990). Such  
505 metallic droplets strongly concentrate and efficiently scavenge highly siderophile  
506 elements (HSEs) from the terrestrial magma ocean and sequester them into the metallic  
507 core. Mantle relative abundances of HSEs resemble the chondrites, leading to the notion  
508 that these elements were delivered during the final ~1% of Earth accretion, after cessation  
509 of core formation (Kimura et al., 1974), now interpreted as accretion after the Moon-  
510 forming giant impact. Isotopic characteristics of the Earth's mantle suggest delivery by  
511 bodies with metallic phases either as undifferentiated metallic grains or as planetesimal  
512 cores (Marchi et al., 2018). However, the terrestrial upper mantle is currently unsaturated

513 in metallic Fe, instead exhibiting a more oxidizing redox state, indicated by higher  
514  $\text{Fe}^{+3}/\text{Fe}^{+2}$  values than those characterizing co-existence with metallic iron. Accordingly,  
515 accreted metals must have been oxidized and dissolved into Earth's mantle, prompting  
516 the question of the nature of the oxidant involved in late accretion. Since the terrestrial  
517 magma ocean crystallized on  $\sim 10^6$  year timescales (Lebrun et al., 2013) while the  
518 leftovers of terrestrial planet accretion were swept up over  $\sim 10^7$ - $10^8$  years (Morbidelli et  
519 al., 2012), late accretion occurred onto a solid silicate Earth with well-developed water  
520 oceans. Possibilities for oxidizing metals delivered during late accretion are: (1) Earth's  
521 fluid envelope, e.g. water in the terrestrial oceans, via the iron-water reaction  
522 ( $\text{Fe}+\text{H}_2\text{O}\rightarrow\text{FeO}+\text{H}_2$ ) followed by hydrodynamic escape of  $\text{H}_2$  (Genda et al., 2017) and  
523 (2) the oxidizing power of Earth's mantle, epitomized by ferric iron ( $2\text{Fe}^{+3}+\text{Fe}^0\rightarrow 3\text{Fe}^{+2}$ ),  
524 lowering the ferric iron abundance towards the modern upper mantle value  
525 ( $\text{Fe}^{+3}/\text{Fe}^{+2}=0.03$ - $0.04$ ) (Canil et al., 1994). Using calculated enrichments in the D/H of  
526 water oceans coexisting with significant ( $\sim 10$ - $100$  bar) early  $\text{H}_2$  atmospheric inventories,  
527 we can limit the extent of the iron-water reaction in oxidizing the metals of late accretion.  
528 Given that the oxidation of metals at that time would consume more than a modern ocean  
529 worth of water via the iron-water reaction (producing  $>30$  bars  $\text{H}_2$ ) and that the Hadean  
530  $\text{H}_2$  inventory was apparently  $<10$ - $20$  bars (Fig. 5), we conclude that the role of the iron-  
531 water reaction during late accretion was subdominant. A more dominant role for this  
532 reaction would have produced deuterium-enriched oceans not observed in the terrestrial  
533 isotopic record. This reasoning suggests that the terrestrial mantle supplied oxidants  
534 during late accretion, a feature that may yield insights into the physics and chemistry of  
535 this early terrestrial process.

536

537 **6. Conclusions**

538 The isotopic composition of the oceans provides a unique constraint for early planetary  
539 evolution. It is widely accepted that most water accreted by the Earth was delivered  
540 before the Moon-forming giant impact and that most water dissolved in the subsequent  
541 magma ocean was excluded from crystallizing minerals and outgassed into the  
542 primordial atmosphere (Elkins-Tanton, 2008; Greenwood et al., 2018). Given that the  
543 residence time of water in the Earth's oceans is comparable to, or greater than, the current  
544 age of the Earth (van Keken et al., 2011), most of the hydrogen in the oceans today is  
545 inferred to be outgassed from the magma ocean, retaining isotopic memory of the  
546 chemical composition of the primordial atmosphere.

547

548 The oxygen fugacity of terrestrial magma ocean outgassing – and therefore the chemical  
549 composition of the primordial atmosphere – has not been independently constrained. By  
550 linking the oxygen fugacity of primordial outgassing to the deuterium content of Earth's  
551 hydrosphere, we articulate new constraints on these critical parameters governing early  
552 Earth evolution. We find that preservation of a carbonaceous chondritic D/H signature  
553 in the terrestrial oceans (to ~10-20%) requires the primordial terrestrial fluid envelope  
554 be hydrogen-poor ( $H_2/H_2O < 0.3$ ), indicating oxidizing conditions ( $\log fO_2 > IW+1$ ) for last  
555 equilibration with the magma ocean. We infer that oxidation of the silicate Earth took  
556 place during the evolution of Earth's final magma ocean, and may require no geological  
557 oxidation processes (e.g. subduction) to be consistent with an oxidized mantle (Trail et  
558 al., 2011) observed in the earliest terrestrial record.

559

560 The inferred absence of massive (>20 bar)  $H_2$  inventories of any origin on the Hadean  
561 Earth constrains the oxidant for terrestrial late accretion. Whereas the likely existence of  
562 early water oceans has previously been taken to imply that the iron-water reaction was

563 responsible for oxidizing the metals delivering HSEs to the early Earth (Genda et al.,  
564 2017), we find that such massive reduction would have disturbed the carbonaceous-  
565 chondrite-like signature of the terrestrial oceans. We therefore infer that oxidants in the  
566 terrestrial mantle (e.g. Fe<sup>+3</sup>) were responsible for oxidative resorption of late-accreting  
567 metals delivered to the Earth. Indeed, the oxidative potential of Earth's modern mantle  
568 (Canil et al., 1994) is comparable to the reducing potential in ~0.5% of an Earth mass of  
569 chondritic late accretion, a feature that may yield insight into this early process.

570

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574

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